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# VANISHING GEODESIC DISTANCE FOR THE RIEMANNIAN METRIC WITH GEODESIC EQUATION THE KDV-EQUATION

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ABSTRACT. The Virasoro-Bott group endowed with the right-invariant  $L^2$ metric (which is a weak Riemannian metric) has the KdV-equation as geodesic equation. We prove that this metric space has vanishing geodesic distance.

## 1. INTRODUCTION

It was found in [11] that a curve in the Virasoro-Bott group is a geodesic for the right invariant  $L^2$ -metric if and only if its right logarithmic derivative is a solution of the Korteweg-de Vries equation, see 2.3. Vanishing geodesic distance for weak Riemannian metrics on infinite dimensional manifolds was first noticed on shape space  $\text{Imm}(S^1, \mathbb{R}^2)/\text{Diff}(S^1)$  for the  $L^2$ -metric in [7, 3.10]. In [8] this result was shown to hold for the general shape space Imm(M, N)/Diff(M) for any compact manifold M and Riemannian manifold N, and also for the right invariant  $L^2$ -metric on each full diffeomorphism group with compact support  $\text{Diff}_c(N)$ . In particular, Burgers' equation is related to the geodesic equation of the right invariant  $L^2$ -metric on  $\text{Diff}(S^1)$  or  $\text{Diff}_c(\mathbb{R})$  and it thus also has vanishing geodesic distance. We even have

**Result.** [8] The weak Riemannian  $L^2$ -metric on each connected component of the total space Imm(M, N) for a compact manifold M and a Riemannian manifold (N, g) has vanishing geodesic distance.

This result is not spelled out in [8] but it follows from there: Given two immersions  $f_0, f_1$  in the same connected component, we first connect their shapes  $f_0(M)$ and  $f_1(M)$  by a curve of length  $< \varepsilon$  in the shape space Imm(M, N) / Diff(M) and take the horizontal lift to get a curve of length  $< \varepsilon$  from  $f_0$  to an immersion  $f_1 \circ \varphi$  in the connected component of the orbit through  $f_1$ . Now we use the induced metric  $f_1^*g$  on M and the right invariant  $L^2$ -metric induced on  $\text{Diff}(M)_0$  to get a curve in Diff(M) of length  $< \varepsilon$  connecting  $\varphi$  with  $\text{Id}_M$ . Evaluating at  $f_1$  we get curve in Imm(M, N) of length  $< \varepsilon$  connecting  $f_1 \circ \varphi$  with  $f_1$ .

In this article we show that the right invariant  $L^2$ -metric on the Virasoro-Bott groups (see 2.1) has vanishing geodesic distance. This might be related to the fact that the Riemannian exponential mapping is not a diffeomorphism near 0, see [2] for Diff $(S^1)$  and [3] for the Virasoro group over  $S^1$ . See [10] for information on conjugate points along geodesics.

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#### 2. The Virasoro-Bott groups

2.1. The Virasoro-Bott groups. Let  $\text{Diff}_{\mathcal{S}}(\mathbb{R})$  be the group of diffeomorphisms of  $\mathbb{R}$  which rapidly fall to the identity. This is a regular Lie group, see [6, 6.4]. The mapping

$$c: \operatorname{Diff}_{\mathcal{S}}(\mathbb{R}) \times \operatorname{Diff}_{\mathcal{S}}(\mathbb{R}) \to \mathbb{R}$$
$$c(\varphi, \psi) := \frac{1}{2} \int \log(\varphi \circ \psi)' d\log \psi' = \frac{1}{2} \int \log(\varphi' \circ \psi) d\log \psi'$$

satisfies  $c(\varphi, \varphi^{-1}) = 0$ ,  $c(\mathrm{Id}, \psi) = 0$ ,  $c(\varphi, \mathrm{Id}) = 0$  and is a smooth group cocycle, called the Bott cocycle:

$$c(\varphi_2,\varphi_3) - c(\varphi_1 \circ \varphi_2,\varphi_3) + c(\varphi_1,\varphi_2 \circ \varphi_3) - c(\varphi_1,\varphi_2) = 0.$$

The corresponding central extension group Vir :=  $\mathbb{R} \times_c \text{Diff}_{\mathcal{S}}(\mathbb{R})$ , called the Virasoro-Bott group, is a trivial  $\mathbb{R}$ -bundle  $\mathbb{R} \times \text{Diff}_{\mathcal{S}}(\mathbb{R})$  that becomes a regular Lie group relative to the operations

$$\begin{pmatrix} \varphi \\ \alpha \end{pmatrix} \begin{pmatrix} \psi \\ \beta \end{pmatrix} = \begin{pmatrix} \varphi \circ \psi \\ \alpha + \beta + c(\varphi, \psi) \end{pmatrix}, \quad \begin{pmatrix} \varphi \\ \alpha \end{pmatrix}^{-1} = \begin{pmatrix} \varphi^{-1} \\ -\alpha \end{pmatrix} \quad \varphi, \psi \in \text{Diff}_{\mathcal{S}}(\mathbb{R}), \ \alpha, \beta \in \mathbb{R}.$$

Other versions of the Virasoro-Bott group are the following:  $\mathbb{R} \times_c \operatorname{Diff}_c(\mathbb{R})$  where  $\operatorname{Diff}_c(\mathbb{R})$  is the group of all diffeomorphisms with compact support, or the periodic case  $\mathbb{R} \times_c \operatorname{Diff}^+(S^1)$ . One can also apply the homomorphism  $\exp(i\alpha)$  to the center and replace it by  $S^1$ . To be specific, we shall treat the most difficult case  $\operatorname{Diff}_{\mathcal{S}}(\mathbb{R})$  in this article. All other cases require only obvious minor changes in the proofs.

2.2. The Virasoro Lie algebra. The Lie algebra of the Virasoro-Bott group  $\mathbb{R} \times_c \text{Diff}_{\mathcal{S}}(\mathbb{R})$  is  $\mathbb{R} \times \mathfrak{X}_{\mathcal{S}}(\mathbb{R})$  (where  $\mathfrak{X}_{\mathcal{S}}(\mathbb{R}) = \mathcal{S}(\mathbb{R})\partial_x$ ) with the Lie bracket

$$\begin{bmatrix} \begin{pmatrix} X \\ a \end{pmatrix}, \begin{pmatrix} Y \\ b \end{bmatrix} = \begin{pmatrix} -[X, Y] \\ \omega(X, Y) \end{pmatrix} = \begin{pmatrix} X'Y - XY' \\ \omega(X, Y) \end{pmatrix}$$

where

$$\omega(X,Y) = \omega(X)Y = \int X'dY' = \int X'Y''dx = \frac{1}{2} \int \det \begin{pmatrix} X' & Y' \\ X'' & Y'' \end{pmatrix} dx,$$

is the *Gelfand-Fuks Lie algebra cocycle*  $\omega : \mathfrak{g} \times \mathfrak{g} \to \mathbb{R}$ , which is a bounded skewsymmetric bilinear mapping satisfying the cocycle condition

$$\omega([X,Y],Z) + \omega([Y,Z],X) + \omega([Z,X],Y) = 0.$$

It is a generator of the 1-dimensional bounded Chevalley cohomology  $H^2(\mathfrak{g}, \mathbb{R})$  for any of the Lie algebras  $\mathfrak{g} = \mathfrak{X}(\mathbb{R}), \ \mathfrak{X}_c(\mathbb{R}), \text{ or } \mathfrak{X}_{\mathcal{S}}(\mathbb{R}) = \mathcal{S}(\mathbb{R})\partial_x$ . The Lie algebra of the Virasoro-Bott Lie group is thus the central extension  $\mathbb{R} \times_{\omega} \mathfrak{X}_{\mathcal{S}}(\mathbb{R})$  induced by this cocycle. We have  $H^2(\mathfrak{X}_c(M), \mathbb{R}) = 0$  for each finite dimensional manifold of dimension  $\geq 2$  (see [4]), which blocks the way to find a higher dimensional analog of the Korteweg-de Vries equation in a way similar to that sketched below.

To complete the description, we add the adjoint action:

$$\operatorname{Ad}\begin{pmatrix}\varphi\\\alpha\end{pmatrix}\begin{pmatrix}Y\\b\end{pmatrix} = \begin{pmatrix}\operatorname{Ad}(\varphi)Y = \varphi_*Y = T\varphi \circ Y \circ \varphi^{-1}\\b + \int S(\varphi)Y \, dx\end{pmatrix}$$

where the *Schwartzian derivative* S is given by

$$S(\varphi) = \left(\frac{\varphi''}{\varphi'}\right)' - \frac{1}{2}\left(\frac{\varphi''}{\varphi'}\right)^2 = \frac{\varphi'''}{\varphi'} - \frac{3}{2}\left(\frac{\varphi''}{\varphi'}\right)^2 = \log(\varphi')'' - \frac{1}{2}(\log(\varphi')')^2$$

which measures the deviation of  $\varphi$  from being a Möbius transformation:

$$S(\varphi) = 0 \iff \varphi(x) = \frac{ax+b}{cx+d} \text{ for } \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL(2,\mathbb{R}).$$

The Schwartzian derivative of a composition and an inverse follow from the action property:

$$S(\varphi \circ \psi) = (S(\varphi) \circ \psi)(\psi')^2 + S(\psi), \quad S(\varphi^{-1}) = -\frac{S(\varphi)}{(\varphi')^2} \circ \varphi^{-1}$$

2.3. The right invariant  $L^2$ -metric and the KdV-equation. We shall use the  $L^2$ -inner product on  $\mathbb{R} \times_{\omega} \mathfrak{X}_{\mathcal{S}}(\mathbb{R})$ :

$$\left\langle \begin{pmatrix} X \\ a \end{pmatrix}, \begin{pmatrix} Y \\ b \end{pmatrix} \right\rangle := \int XY \, dx + ab.$$

We use the induced right invariant weak Riemannian metric on the Virasoro group.

According to [1], see [9] for a proof in the notation and setup used here, a curve  $t \mapsto \begin{pmatrix} \varphi(t, \ \alpha(t) \end{pmatrix}$  in the Virasoro-Bott group is a geodesic if and only if

$$\begin{pmatrix} u_t \\ a_t \end{pmatrix} = -\operatorname{ad} \begin{pmatrix} u \\ a \end{pmatrix}^{\top} \begin{pmatrix} u \\ a \end{pmatrix} = \begin{pmatrix} -3u_x u - au_{xxx} \\ 0 \end{pmatrix} \text{ where}$$

$$\begin{pmatrix} u(t) \\ a(t) \end{pmatrix} = \partial_s \begin{pmatrix} \varphi(s) \\ \alpha(s) \end{pmatrix} \cdot \begin{pmatrix} \varphi(t)^{-1} \\ -\alpha(t) \end{pmatrix} \Big|_{s=t} = \partial_s \begin{pmatrix} \varphi(s) \circ \varphi(t)^{-1} \\ \alpha(s) - \alpha(t) + c(\varphi(s), \varphi(t)^{-1}) \end{pmatrix} \Big|_{s=t},$$

$$\begin{pmatrix} u \\ a \end{pmatrix} = \begin{pmatrix} \varphi_t \circ \varphi^{-1} \\ \alpha_t - \int \frac{\varphi_{tx}\varphi_{xx}}{2\varphi_x^2} dx \end{pmatrix},$$

since we have

$$\begin{aligned} 2\partial_s c(\varphi(s),\varphi(t)^{-1})|_{s=t} &= \partial_s \int \log(\varphi(s)' \circ \varphi(t)^{-1}) d\log((\varphi(t)^{-1})')|_{s=t} \\ &= \int \frac{\varphi_t(t)' \circ \varphi(t)^{-1}}{\varphi(t)' \circ \varphi(t)^{-1}} \left( -\frac{\varphi(t)'' \circ \varphi(t)^{-1}}{(\varphi(t)' \circ \varphi(t)^{-1})^2} \right) dx \\ &= -\int \left( \frac{\varphi_t' \varphi''}{(\varphi')^2} \right) (t) dy = -\int \left( \frac{\varphi_{tx} \varphi_{xx}}{\varphi_x^2} \right) (t) dx. \end{aligned}$$

Thus a is a constant in time and the geodesic equation is hence the Kortewegde Vries equation

$$u_t + 3u_x u + au_{xxx} = 0$$

with its natural companions

$$\varphi_t = u \circ \varphi, \qquad \alpha_t = a + \int \frac{\varphi_{tx} \varphi_{xx}}{2\varphi_x^2} dx.$$

To be complete, we add the invariant momentum mapping J with values in the Virasoro algebra (via the weak Riemannian metric). We need the transpose of the adjoint action:

$$\left\langle \operatorname{Ad} \begin{pmatrix} \varphi \\ \alpha \end{pmatrix}^{\top} \begin{pmatrix} Y \\ b \end{pmatrix}, \begin{pmatrix} Z \\ c \end{pmatrix} \right\rangle = \left\langle \begin{pmatrix} Y \\ b \end{pmatrix}, \operatorname{Ad} \begin{pmatrix} \varphi \\ \alpha \end{pmatrix} \begin{pmatrix} Z \\ c \end{pmatrix} \right\rangle$$
$$= \left\langle \begin{pmatrix} Y \\ b \end{pmatrix}, \begin{pmatrix} \varphi_* Z \\ c + \int S(\varphi) Z \, dx \end{pmatrix} \right\rangle$$

$$= \int Y((\varphi' \circ \varphi^{-1})(Z \circ \varphi^{-1}) \, dx + bc + \int bS(\varphi)Z \, dx$$
$$= \int ((Y \circ \varphi)(\varphi')^2 + bS(\varphi))Z \, dx + bc$$

Thus, the invariant momentum mapping is given by

$$J\left(\begin{pmatrix}\varphi\\\alpha\end{pmatrix},\begin{pmatrix}Y\\b\end{pmatrix}\right) = \operatorname{Ad}\begin{pmatrix}\varphi\\\alpha\end{pmatrix}^{\top}\begin{pmatrix}Y\\b\end{pmatrix} = \begin{pmatrix}(Y\circ\varphi)(\varphi')^2 + bS(\varphi)\\b\end{pmatrix}.$$

Along a geodesic  $t \mapsto g(t, \ ) = \begin{pmatrix} \varphi(t, \ ) \\ \alpha(t) \end{pmatrix}$ , the momentum

$$J\left(\binom{\varphi}{\alpha}, \binom{u = \varphi_t \circ \varphi^{-1}}{a}\right) = \binom{(u \circ \varphi)\varphi_x^2 + aS(\varphi)}{a} = \binom{\varphi_t \varphi_x^2 + aS(\varphi)}{a}$$

is constant in t.

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## 2.4. Lifting curves to the Virasoro-Bott group. We consider the extension

$$\mathbb{R} \xrightarrow{i} \mathbb{R} \times_c \operatorname{Diff}_{\mathcal{S}}(\mathbb{R}) \xrightarrow{p} \operatorname{Diff}_{\mathcal{S}}(\mathbb{R}).$$

Then p is a Riemannian submersion for the right invariant  $L^2$ -metric on  $\text{Diff}_{\mathcal{S}}(\mathbb{R})$ , i.e., Tp is an isometry on the orthogonal complements of the fibers. These complements are not integrable; in fact, the curvature of the corresponding principal connection is given by the Gelfand-Fuks cocycle. For any curve  $\varphi(t)$  in  $\text{Diff}_{\mathcal{S}}(\mathbb{R})$ its horizontal lift is given by

$$\begin{pmatrix} \varphi(t) \\ a(t) = a(0) - \int_0^t \int \frac{\varphi_{tx}\varphi_{xx}}{\varphi_x^2} \, dx \, dt \end{pmatrix}$$

since the right translation to (Id, 0) of its velocity should have zero vertical component, see 2.3. The horizontal lift has the same length and energy as  $\varphi$ .

### 3. VANISHING OF THE GEODESIC DISTANCE

3.1. **Theorem.** On all Virasoro-Bott groups mentioned in 2.1 geodesic distance for the right invariant  $L^2$ -metric vanishes.

The rest of this section is devoted to the proof of theorem 3.1 for the most difficult case  $\mathbb{R} \times_c \text{Diff}_{\mathcal{S}}(\mathbb{R})$ .

3.2. **Proposition.** Any two diffeomorphisms in  $\text{Diff}_{\mathcal{S}}(\mathbb{R})$  can be connected by a path with arbitrarily short length for the right invariant  $L^2$ -metric.

In [8] for  $\operatorname{Diff}_c(\mathbb{R})$  it was first shown that there exists one non-trivial diffeomorphism which can be connected to Id with arbitrarily small length. Then, it was shown that the diffeomorphisms with this property form a normal subgroup. Since  $\operatorname{Diff}_c(\mathbb{R})$  is a simple group this concluded the proof. But  $\operatorname{Diff}_{\mathcal{S}}(\mathbb{R})$  is not a simple group since  $\operatorname{Diff}_c(\mathbb{R})$  is a normal subgroup. So, we have to elaborate on the proof of [8] as follows.

**Proof.** We show that any rapidly decreasing diffeomorphism can be connected to the identity by an arbitrarily short path. We will write this diffeomorphism as  $\operatorname{Id} + g$ , where  $g \in \mathcal{S}(\mathbb{R})$  is a rapidly decreasing function with g' > -1. For  $\lambda = 1 - \varepsilon < 1$  we define

 $\varphi(t,x) = x + \max(0,\min(t - \lambda x, g(x))) - \max(0,\min(t + \lambda x, -g(x))).$ 

This is a (non-smooth) path defined for  $t \in (-\infty, \infty)$  connecting the identity in  $\operatorname{Diff}_{\mathcal{S}}(\mathbb{R})$  with the diffeomorphism  $(\operatorname{Id} + g)$ . We define  $\psi(t, x) = \varphi(\operatorname{tan}(t), x) \star G_{\varepsilon}(t, x)$ , where  $G_{\varepsilon}(t, x) = \frac{1}{\varepsilon^2}G_1(\frac{t}{\varepsilon}, \frac{x}{\varepsilon})$  is a smoothing kernel with  $\operatorname{supp}(G_{\varepsilon}) \subseteq B_{\varepsilon}(0)$ and  $\iint G_{\varepsilon} \, dx \, dt = 1$ . Thus  $\psi$  is a smooth path defined on the finite interval  $-\frac{\pi}{2} < t < \frac{\pi}{2}$  connecting the identity in  $\operatorname{Diff}_{\mathcal{S}}(\mathbb{R})$  with a diffeomorphism arbitrarily close to  $(\operatorname{Id} + g)$  for  $\varepsilon$  small. (Compare figure 1 for an illustration.)

The  $L^2$ -energy of  $\psi$  is

$$E(\psi) = \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} \int_{\mathbb{R}} (\psi_t \circ \psi^{-1})^2 \, \mathrm{d}x \, \mathrm{d}t = \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} \int_{\mathbb{R}} \psi_t^2 \psi_x \, \mathrm{d}x \, \mathrm{d}t$$

where  $\psi^{-1}(t, x)$  stands for  $\psi(t, -)^{-1}(x)$ . We have

$$\partial_a \max(0, \min(a, b)) = \mathbb{1}_{0 \le a \le b}, \qquad \partial_b \max(0, \min(a, b)) = \mathbb{1}_{0 \le b \le a},$$

and therefore

$$\begin{split} \psi_x(t,x) &= \varphi_x(\tan(t),x) \star G_{\varepsilon} \\ &= \left(1 - \lambda \mathbb{1}_{0 \leq \tan(t) - \lambda x \leq g(x)} + g'(x) \mathbb{1}_{0 \leq g(x) \leq \tan(t) - \lambda x} \right. \\ &\quad - \lambda \mathbb{1}_{0 \leq \tan(t) + \lambda x \leq -g(x)} + g'(x) \mathbb{1}_{0 \leq -g(x) \leq \tan(t) + \lambda x} \right) \star G_{\varepsilon}, \\ \psi_t(t,x) &= \left((1 + \tan(t)^2)\varphi_t(\tan(t),x)\right) \star G_{\varepsilon} \\ &= \left((1 + \tan(t)^2)(\mathbb{1}_{0 \leq \tan(t) - \lambda x \leq g(x)} - \mathbb{1}_{0 \leq \tan(t) + \lambda x \leq -g(x)})\right) \star G_{\varepsilon}. \end{split}$$

Note that these functions have disjoint support when  $\varepsilon = 0, \lambda = 1 - \varepsilon = 1$ . **Claim.** The mappings  $\varepsilon \mapsto \psi_t$  and  $\varepsilon \mapsto (\psi_x - 1)$  are continuous into each  $L^p$  with p even. (The proofs are simpler when p is even because there are no absolute values to be taken care of.) To prove the claim, we calculate

$$\begin{split} \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} &\int_{\mathbb{R}} \left( (1 + \tan(t)^2) \varphi_t(\tan(t), x) \right)^p \mathrm{d}x \, \mathrm{d}t = \iint_{\mathbb{R}^2} \varphi_t(t, x)^p (1 + t^2)^{p-1} \, \mathrm{d}x \, \mathrm{d}t \\ &= \iint_{\mathbb{R}^2} (\mathbb{1}_{0 \le t - \lambda x \le g(x)} + \mathbb{1}_{0 \le t + \lambda x \le -g(x)}) (1 + t^2)^{p-1} \, \mathrm{d}x \, \mathrm{d}t \\ &= \int_{g(x) \ge 0} \int_{\lambda x}^{\lambda x + g(x)} (1 + t^2)^{p-1} \, \mathrm{d}t \, \mathrm{d}x + \int_{g(x) < 0} \int_{\lambda x + g(x)}^{\lambda x} (1 + t^2)^{p-1} \, \mathrm{d}t \, \mathrm{d}x \\ &= \int_{\mathbb{R}} \left| F(t) \right|_{t = \lambda x}^{t = \lambda x + g(x)} \right| \, \mathrm{d}x = \int_{\mathbb{R}} |F(\lambda x + g(x)) - F(\lambda x)| \, \mathrm{d}x, \end{split}$$

where  $F(\lambda x + g(x)) - F(\lambda x)$  is a polynomial without constant term in g(x) with coefficients also powers of  $\lambda x$ . Integrals of the form  $\int_{\mathbb{R}} |(\lambda x)^{k_1}g(x)^{k_2}| dx$  with  $k_1 \geq 0, k_2 > 0$  are finite and continuous in  $\lambda = 1 - \varepsilon$  since g is rapidly decreasing. This shows that  $||(1 + \tan(t)^2)\varphi_t(\tan(t), x)||_p$  depends continuously on  $\varepsilon$ . Furthermore the sequence  $(1 + \tan(t)^2)\varphi_t(\tan(t), x)$  converges almost everywhere for  $\varepsilon \to 0$ , thus it also converges in measure. By the theorem of Vitali, this implies convergence in  $L^p$ , see for example [12, theorem 16.6]. Convolution with  $G_{\varepsilon}$  acts as approximate unit in each  $L^p$ , which proves the claim for  $\psi_t$ . For  $\psi_x - 1$  it follows similarly.

The above claim implies that

$$E(\psi) = \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} \int_{\mathbb{R}} \psi_t^2 \psi_x \, \mathrm{d}x \, \mathrm{d}t = \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} \int_{\mathbb{R}} \psi_t^2 (\psi_x - 1) \, \mathrm{d}x \, \mathrm{d}t + \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} \int_{\mathbb{R}} \psi_t^2 \, \mathrm{d}x \, \mathrm{d}t$$

viewed as a mapping on  $L^4 \times L^4 \times L^2$  (first summand) and on  $L^2 \times L^2$  (second summand) is continuous in  $\varepsilon$ . It also vanishes at  $\varepsilon = 0$  since then  $\psi_x$  and  $\psi_t$  have disjoint support. The Cauchy-Schwarz inequality  $L(\psi)^2 < \pi E(\psi)$  implies that  $L(\psi)$  goes to zero as well. Ultimately,  $\psi(\frac{\pi}{2}) = (\mathrm{Id} + g) \star G_{\varepsilon}$  is arbitrarily close to  $\mathrm{Id} + g$ .

3.3. Lemma. For any  $a \in \mathbb{R}$  there exists an arbitrarily short path connecting  $\binom{\mathrm{Id}}{0}$  and  $\binom{\mathrm{Id}}{a}$ , *i.e.*,  $\operatorname{dist}_{\operatorname{Vir}}^{L^2}\left(\binom{\mathrm{Id}}{0}, \binom{\mathrm{Id}}{a}\right) = 0$ .

**Proof.** The aim of the following argument is to construct a family of paths in the diffeomorphism group, parametrized by  $\varepsilon$ , with the following properties: all paths in the family start and end at the identity and their length in the diffeomorphism group with respect to the  $L^2$  metric tends to 0 as  $\varepsilon \to 0$ . By letting  $\varepsilon$  be time-dependent, we are able to control the endpoint a(T) of the horizontal lift for certain diffeomorphisms.

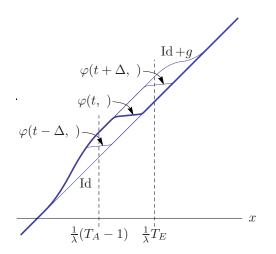


FIGURE 1. The path  $\varphi(t, \cdot)$  defined in 3.3 connecting Id to Id +g, plotted at  $t - \Delta < t < t + \Delta$ . Between the dashed lines,  $g \equiv 1$  is constant.

We consider the function

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(1)  

$$f(z, a, \varepsilon) = \max(0, \min(z, a)) \star G_{\varepsilon}(z)G_{\varepsilon}(a)$$

$$= \iint \max(0, \min(z - \overline{z}, a - \overline{a}))G_{\varepsilon}(\overline{z})G_{\varepsilon}(\overline{a}) \, \mathrm{d}\overline{z} \, \mathrm{d}\overline{a}$$

$$= \iint \max(0, \min(z - \varepsilon \overline{z}, a - \varepsilon \overline{a}))G_{1}(\overline{z})G_{1}(\overline{a}) \, \mathrm{d}\overline{z} \, \mathrm{d}\overline{a}$$

$$= \varepsilon f(\frac{z}{\varepsilon}, \frac{a}{\varepsilon}, 1)$$

where  $G_{\varepsilon}(z) = \frac{1}{\varepsilon}G_1(\frac{z}{\varepsilon})$  is a function with  $\operatorname{supp}(G_{\varepsilon}) \subseteq [-\varepsilon, \varepsilon]$  and  $\int G_{\varepsilon} dx = 1$ . Furthermore, let  $g: \mathbb{R} \to [0, 1]$  be a function with compact support contained in  $\mathbb{R}_{>0}$  and g' > -1, so that x + g(x) is a diffeomorphism. For  $0 < \lambda < 1$  and  $t \in [0, T]$  let

$$\varphi(t, x) = x + f(t - \lambda x, g(x), \varepsilon(t))$$

be the path going away from the identity (since  $\operatorname{supp}(g) \subset \mathbb{R}_{>0}$ , see also figure 1). For given  $\varepsilon_0 > 0$ , let

$$\psi(t, x) = x + f(T - t - \lambda x, g(x), \varepsilon_0)$$

the path leading back again. The only difference to [8] is that the parameter  $\varepsilon$  may vary along the path.

We shall need some derivatives of  $\varphi$  and f:

$$\begin{split} \varphi_t(t,x) &= f_z(t-\lambda x,g(x),\varepsilon(t)) + \dot{\varepsilon}(t)f_\varepsilon(t-\lambda x,g(x),\varepsilon(t)) \\ \varphi_x(t,x) &= 1-\lambda f_z(t-\lambda x,g(x),\varepsilon(t)) + f_a(t-\lambda x,g(x),\varepsilon(t))g'(x) \\ f_z(z,a,\varepsilon) &= \int_{-\infty}^z \int_{-\infty}^{a-z} G_\varepsilon(w)G_\varepsilon(w+b)\,\mathrm{d}b\,\mathrm{d}w \\ f_a(z,a,\varepsilon) &= \int_{-\infty}^a \int_{-\infty}^{z-a} G_\varepsilon(w)G_\varepsilon(w+b)\,\mathrm{d}b\,\mathrm{d}w \\ f_\varepsilon(z,a,\varepsilon) &= \frac{1}{\varepsilon} \Big( f(z,a,\varepsilon) - zf_z(z,a,\varepsilon) - af_a(z,a,\varepsilon) \Big) \\ f_{zz}(z,a,\varepsilon) &= G_\varepsilon(z) \int_{-\infty}^a G_\varepsilon(b)\,\mathrm{d}b - \int_{-\infty}^z G_\varepsilon(w)G_\varepsilon(w-(z-a))\,\mathrm{d}w \end{split}$$

**Claim 1.** The path  $\varphi$  followed by  $\psi$  still has arbitrarily small length for the  $L^2$ -metric.

We are working with a fixed time interval [0, 2T]. Thus arbitrarily small length is equivalent to arbitrarily small energy. The energy is given by

(2) 
$$\iint \varphi_t^2 \varphi_x \, \mathrm{d}x \, \mathrm{d}t = \iint (f_z + \dot{\varepsilon} f_\varepsilon)^2 (1 - \lambda f_z + f_a g') \, \mathrm{d}x \, \mathrm{d}t$$

Looking at the formula for  $f_{\varepsilon}$  we see that  $\varepsilon f_{\varepsilon}$  is bounded on a domain with bounded a. Thus  $\|\dot{\varepsilon}f_{\varepsilon}\|_{\infty} \to 0$  can be achieved by choosing  $\varepsilon$ , such that  $|\dot{\varepsilon}| \leq C\varepsilon^{3/2}$ . We will see later that this is possible. Inspecting  $\varphi_t(t, x)$  and looking at the formulas for  $f_z$  and f we see that for  $t - \lambda x < -\varepsilon(t)$  and for  $t - \lambda x - g(x) > 2\varepsilon(t)$  we have  $\varphi_t(t, x) = 0$ . Thus the domain of integration is contained in the compact set

$$[0,T] \times [-\frac{T + \|g\|_{\infty} + 2\|\varepsilon\|_{\infty}}{\lambda}, \frac{T + \|\varepsilon\|_{\infty}}{\lambda}].$$

Therefore, it is enough to show that the  $L^{\infty}$ -norm of the integrand in (2) goes to zero as  $\|\varepsilon\|_{\infty}$  goes to zero. For all terms involving  $\dot{\varepsilon}f_{\varepsilon}$  this is true by the above assumption since  $(1 - \lambda f_z + f_a g')$  and  $\varepsilon f_{\varepsilon}$  are bounded. For the remaining parts  $f_z^2(1 - \lambda f_z)$  and  $f_z^2 f_a g'$  we follow the argumentation of [8]. For t fixed and  $\lambda$  close to 1, the function  $1 - \lambda f_z$ , when restricted to the support of  $f_z$ , is bigger than  $\varepsilon(t)$ only on an interval of length  $O(\varepsilon(t))$ . Hence we have

$$\int_0^T \int_{\mathbb{R}} f_z^2 (1 - \lambda f_z) \, \mathrm{d}x \, \mathrm{d}t \le \|f_z\|_\infty^2 \int_0^T \int_{\mathbb{R}} (1 - \lambda f_z) \, \mathrm{d}x \, \mathrm{d}t = O(\|\varepsilon\|_\infty)$$

For the last part, we note that the support of  $f_z^2 f_a$  is contained in the set  $|g(x) - (t - \lambda x)| \le 2\varepsilon$ . Now we define  $x_0 < x_1$  by  $g(x_0) + \lambda x_0 = T - 2||\varepsilon||_{\infty}$  and  $g(x_1) + \lambda x_1 = T + 2||\varepsilon||_{\infty}$ . Then

$$\int_0^T \int_{\mathbb{R}} f_z^2 f_a g' \, \mathrm{d}x \, \mathrm{d}t \le T \|f_z\|_\infty^2 \|f_a\|_\infty \int_{\mathrm{supp}(f_z^2 f_a)} g' \, \mathrm{d}x$$

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$$= T(g(x_1) - g(x_0)) \le 4T \|\varepsilon\|_{\infty}.$$

The estimate for  $\psi$  is similar and easier. This proves claim 1. Claim 2. For every  $a \in \mathbb{R}$  and  $\delta > 0$  we may choose  $\varepsilon(t)$  with  $\|\varepsilon\|_{\infty} < \delta$  such that

$$\int_0^T \int_{\mathbb{R}} \frac{\varphi_{tx} \varphi_{xx}}{\varphi_x^2} \, \mathrm{d}x \, \mathrm{d}t + \int_0^T \int_{\mathbb{R}} \frac{\psi_{tx} \psi_{xx}}{\psi_x^2} \, \mathrm{d}x \, \mathrm{d}t = a.$$

We will subject  $\varepsilon$  and g to several assumptions. First, we partition the interval [0,T] equidistantly into  $0 < T_A < T_E < T$  and the (t,x)-domain into two parts, namely  $A_1 = ([0,T_A] \cup [T_E,T]) \times \mathbb{R}$  and  $A_2 = [T_A,T_E] \times \mathbb{R}$ . We want  $g(x) \equiv 1$  on a neighborhood of the interval  $[\frac{1}{\lambda}(T_A-1), \frac{1}{\lambda}T_E]$ . We choose  $\varepsilon(t)$  to be constant  $\varepsilon(t) \equiv \varepsilon_0$  on  $[0,T_A] \cup [T_E,T]$  and to be symmetric in the sense, that  $\varepsilon(t) = \varepsilon(T-t)$ . In addition, we want  $\varepsilon(t)$  small enough, such that  $g(x) \equiv 1$  on  $[\frac{1}{\lambda}(T_A-1-2\varepsilon(t)), \frac{1}{\lambda}(T_E+\varepsilon(t))]$ .

On  $A_1$  we have  $\varepsilon(t) \equiv \varepsilon_0$ . This implies  $\psi_{tx}(t,x) = -\varphi_{tx}(T-t,x)$ ,  $\psi_x(t,x) = \varphi_x(T-t,x)$  and  $\psi_{xx}(t,x) = \varphi_{xx}(T-t,x)$ . Hence

$$\iint_{A_1} \frac{\varphi_{tx}\varphi_{xx}}{\varphi_x^2} \, \mathrm{d}x \, \mathrm{d}t + \iint_{A_1} \frac{\psi_{tx}\psi_{xx}}{\psi_x^2} \, \mathrm{d}x \, \mathrm{d}t = 0.$$

Let  $A_2 = [T_A, T_E] \times \mathbb{R}$  be the region, where  $\varepsilon(t)$  is not constant. In the interior, where

$$\begin{aligned} & -\varepsilon(t) & < t - \lambda x < g(x) + 2\varepsilon(t) \\ & t - g(x) - 2\varepsilon(t) & < \lambda x < t + \varepsilon(t) \end{aligned}$$

we have by assumption  $g(x) \equiv 1$ . Therefore, one has in this region:

$$\begin{aligned} \varphi_x(t,x) &= -\lambda f_z(t - \lambda x, 1, \varepsilon(t)) \\ \varphi_{xx}(t,x) &= \lambda^2 f_{zz}(t - \lambda x, 1, \varepsilon(t)) \\ \varphi_{tx}(t,x) &= -\lambda f_{zz}(t - \lambda x, 1, \varepsilon(t)) - \lambda f_{\varepsilon z}(t - \lambda x, 1, \varepsilon(t)) \dot{\varepsilon}(t) \end{aligned}$$

We divide the integral over  $A_2$  into two symmetric parts

$$\int_{T_A}^{T/2} \int_{\frac{1}{\lambda}(t-1-2\varepsilon(t))}^{\frac{1}{\lambda}(t+\varepsilon(t))} \frac{\varphi_{tx}\varphi_{xx}}{\varphi_x^2} \, \mathrm{d}x \, \mathrm{d}t + \int_{T/2}^{T_E} \int_{\frac{1}{\lambda}(t-1-2\varepsilon(t))}^{\frac{1}{\lambda}(t+\varepsilon(t))} \frac{\varphi_{tx}\varphi_{xx}}{\varphi_x^2} \, \mathrm{d}x \, \mathrm{d}t$$

and apply the following variable substitution to the second integral

$$\widetilde{t} = T - t, \quad \widetilde{x} = x + \frac{1}{\lambda}(\widetilde{t} - t).$$

Thus  $\tilde{t} - \lambda \tilde{x} = t - \lambda x$ . Together with  $\varepsilon(t) = \varepsilon(\tilde{t})$  this implies

$$\varphi_x(t,x) = \varphi_x(\widetilde{t},\widetilde{x}), \quad \varphi_{xx}(t,x) = \varphi_{xx}(\widetilde{t},\widetilde{x}).$$

Since  $\dot{\varepsilon}(t) = -\dot{\varepsilon}(\tilde{t})$  changes sign, the term containing  $\dot{\varepsilon}(t)$  cancels out and leaves only

$$\varphi_{tx}(t,x) + \varphi_{tx}(t,\widetilde{x}) = -2\lambda f_{zz}(t-\lambda x, 1,\varepsilon(t)).$$

A simple calculation shows that the integration limits transform

$$\int_{T/2}^{T_E} \int_{\frac{1}{\lambda}(t-1-2\varepsilon)}^{\frac{1}{\lambda}(t+\varepsilon)} \frac{\varphi_{tx}\varphi_{xx}}{\varphi_x^2} \, \mathrm{d}x \, \mathrm{d}t = \int_{T_A}^{T/2} \int_{\frac{1}{\lambda}(\tilde{t}+\varepsilon)}^{\frac{1}{\lambda}(\tilde{t}+\varepsilon)} \frac{\varphi_{tx}\varphi_{xx}}{\varphi_x^2} \, \mathrm{d}\tilde{x} \, \mathrm{d}\tilde{t}$$

to those of the first integral. Therefore, the sum of the integrals gives

$$\iint_{A_2} \frac{\varphi_{tx}\varphi_{xx}}{\varphi_x^2} \,\mathrm{d}x \,\mathrm{d}t = -2\lambda^3 \int_{T_A}^{T/2} \int_{\frac{1}{\lambda}(t+\varepsilon(t))}^{\frac{1}{\lambda}(t+\varepsilon(t))} \frac{f_{zz}(t-\lambda x, 1, \varepsilon(t))^2}{\left(1-\lambda f_z(t-\lambda x, 1, \varepsilon(t))\right)^2} \,\mathrm{d}x \,\mathrm{d}t.$$

From formula (1) we see:

$$f_z(z, a, \varepsilon) = f_z(\frac{z}{\varepsilon}, \frac{a}{\varepsilon}, 1), \qquad f_{zz}(z, a, \varepsilon) = \frac{1}{\varepsilon} f_{zz}(\frac{z}{\varepsilon}, \frac{a}{\varepsilon}, 1).$$

We can use this to rewrite the above integral:

$$\begin{split} \iint_{A_2} \frac{\varphi_{tx}\varphi_{xx}}{\varphi_x^2} \, \mathrm{d}x \, \mathrm{d}t &= -2\lambda^3 \int_{T_A}^{T/2} \int_{\frac{1}{\lambda}(t-1-2\varepsilon(t))}^{\frac{1}{\lambda}(t+\varepsilon(t))} \frac{f_{zz}(t-\lambda x,1,\varepsilon(t))^2}{\left(1-\lambda f_z(t-\lambda x,1,\varepsilon(t))\right)^2} \, \mathrm{d}x \, \mathrm{d}t \\ &= -2\lambda^2 \int_{T_A}^{T/2} \int_{-\varepsilon(t)}^{2\varepsilon(t)+1} \frac{f_{zz}(z,1,\varepsilon(t))^2}{\left(1-\lambda f_z(z,1,\varepsilon(t))\right)^2} \, \mathrm{d}z \, \mathrm{d}t \\ &= -2\lambda^2 \int_{T_A}^{T/2} \int_{-\varepsilon(t)}^{2\varepsilon(t)+1} \frac{1}{\varepsilon(t)^2} \frac{f_{zz}(\frac{z}{\varepsilon(t)},\frac{1}{\varepsilon(t)},1)^2}{\left(1-\lambda f_z(\frac{z}{\varepsilon(t)},\frac{1}{\varepsilon(t)},1)\right)^2} \, \mathrm{d}z \, \mathrm{d}t \\ &= -2\lambda^2 \int_{T_A}^{T/2} \int_{-\varepsilon(t)}^{2\varepsilon(t)+1} \frac{1}{\varepsilon(t)} \frac{f_{zz}(z,\frac{1}{\varepsilon(t)},\frac{1}{\varepsilon(t)},1)^2}{\left(1-\lambda f_z(z,\frac{1}{\varepsilon(t)},1)\right)^2} \, \mathrm{d}z \, \mathrm{d}t \end{split}$$

Looking at the formula for  $f_{zz}$ 

$$f_{zz}(z, \frac{1}{\varepsilon}, 1) = G_1(z) - \int_{-\infty}^{z} G_1(w) G_1(w - (z - \frac{1}{\varepsilon})) dw$$

we see that  $f_{zz}(z, \frac{1}{\varepsilon}, 1)$  is non-zero only on the intervals |z| < 1 and  $|z - \frac{1}{\varepsilon}| < 2$ . For small  $\varepsilon$ , these are two disjoint regions. Therefore, the above integral equals

$$\iint_{A_2} \frac{\varphi_{tx}\varphi_{xx}}{\varphi_x^2} \, \mathrm{d}x \, \mathrm{d}t = -2\lambda^2 \int_{T_A}^{T/2} \frac{1}{\varepsilon(t)} \int_{-1}^1 \frac{f_{zz}(z, \frac{1}{\varepsilon(t)}, 1)^2}{\left(1 - \lambda f_z(z, \frac{1}{\varepsilon(t)}, 1)\right)^2} \, \mathrm{d}z \, \mathrm{d}t - 2\lambda^2 \int_{T_A}^{T/2} \frac{1}{\varepsilon(t)} \int_{-2}^2 \frac{f_{zz}(z + \frac{1}{\varepsilon(t)}, \frac{1}{\varepsilon(t)}, 1)^2}{\left(1 - \lambda f_z(z + \frac{1}{\varepsilon(t)}, \frac{1}{\varepsilon(t)}, 1)\right)^2} \, \mathrm{d}z \, \mathrm{d}t$$

For z bounded and sufficiently small  $\varepsilon(t)$ , the functions under the integral do not depend on  $\varepsilon(t)$  any more as can be seen from the definitions of  $f_z$  and  $f_{zz}$ . Thus

$$I = \lambda^2 \int_{-1}^1 \frac{f_{zz}(z, \frac{1}{\varepsilon(t)}, 1)^2}{\left(1 - \lambda f_z(z, \frac{1}{\varepsilon(t)}, 1)\right)^2} \,\mathrm{d}z + \lambda^2 \int_{-2}^2 \frac{f_{zz}(z + \frac{1}{\varepsilon(t)}, \frac{1}{\varepsilon(t)}, 1)^2}{\left(1 - \lambda f_z(z + \frac{1}{\varepsilon(t)}, \frac{1}{\varepsilon(t)}, 1)\right)^2} \,\mathrm{d}z,$$

is independent of t and we have

$$\iint_{A_2} \frac{\varphi_{tx} \varphi_{xx}}{\varphi_x^2} \, \mathrm{d}x \, \mathrm{d}t = -I \int_{T_A}^{T_E} \frac{1}{\varepsilon(t)} \, \mathrm{d}t.$$

The same calculations can be repeated for the return path  $\psi$ , where  $\varepsilon \equiv \varepsilon_0$  is constant in time:

$$\iint_{A_2} \frac{\psi_{tx}\psi_{xx}}{\psi_x^2} \,\mathrm{d}x \,\mathrm{d}t = I \int_{T_A}^{T_E} \frac{1}{\varepsilon_0} \,\mathrm{d}t$$

Note that the sign is positive now, which comes from the t-derivative. Putting everything together gives us

$$a = \iint \frac{\varphi_{tx}\varphi_{xx}}{\varphi_x^2} + \frac{\psi_{tx}\psi_{xx}}{\psi_x^2} \, \mathrm{d}x \, \mathrm{d}t = I \int_{T_A}^{T_E} \left(\frac{1}{\varepsilon_0} - \frac{1}{\varepsilon(t)}\right) \mathrm{d}t$$

Let  $\varepsilon(t) = \varepsilon_0 + \varepsilon_1 \varepsilon_0^{3/2} b(t)$  where b(t) is a bump function with height 1 and  $\varepsilon_1$  is a small constant. Note that  $\varepsilon(t)$  satisfies  $|\dot{\varepsilon}| \leq ||\dot{b}||_{\infty} \varepsilon_1 \varepsilon_0^{3/2}$ . Choosing  $\varepsilon_0$  and  $\varepsilon_1$  small independently we may produce any  $a \in \mathbb{R}$ .

**Proof of Theorem 3.1.** Let  $(\varphi, a) \in \mathbb{R} \times_c \text{Diff}_{\mathcal{S}}(\mathbb{R})$ . By proposition 3.2 we get a smooth family  $\varphi(\delta, t, x)$  for  $\delta > 0$  and  $t \in [0, 1]$  such that  $\varphi(\delta, t, \cdot) \in \text{Diff}_{\mathcal{S}}(\mathbb{R})$ ,  $\varphi(\delta, 0, \cdot) = \text{Id}_{\mathbb{R}}, \varphi(\delta, 1, \cdot) = \varphi$ , and such that the length of  $t \mapsto \varphi(\delta, t, \cdot)$  is  $< \delta$ .

Using 2.4 consider the horizontal lift  $(\varphi(\delta, t, ), a(\delta, t)) \in \text{Diff}_{\mathcal{S}}(\mathbb{R})$  of this family which connects  $\binom{\text{Id}}{0}$  with  $\binom{\varphi}{a(\delta,1)}$  for each  $\delta > 0$  and has length  $< \delta$ . But one can see from the proof of lemma 3.3 that  $a(\delta, 1)$  becomes unbounded for  $\delta \to 0$ .

Using lemma 3.3 we can find a horizontal path  $t \mapsto \begin{pmatrix} \psi(\delta,t, \cdot) \\ b(\delta,t) \end{pmatrix}$  for  $t \in [0,1]$  in the Virasoro group of length  $<\delta$  connecting  $\begin{pmatrix} \mathrm{Id} \\ 0 \end{pmatrix}$  with  $\begin{pmatrix} \mathrm{Id} \\ a-a(\delta,1) \end{pmatrix}$ . Then the curve  $t \mapsto \begin{pmatrix} \psi(\delta,t, \cdot) \\ b(\delta,t) \end{pmatrix} \cdot \begin{pmatrix} \varphi \\ a(\delta,1) \end{pmatrix} = \begin{pmatrix} \psi(\delta,t) \circ \varphi \\ b(\delta,t) + a(\delta,1) + c(\psi(\delta,t),\varphi) \end{pmatrix}$  connects  $\begin{pmatrix} \varphi \\ a(\delta,1) \end{pmatrix} = \begin{pmatrix} \mathrm{Id} \\ 0 \end{pmatrix} \cdot \begin{pmatrix} \varphi \\ a(\delta,1) \end{pmatrix}$  with  $\begin{pmatrix} \varphi \\ a \end{pmatrix} = \begin{pmatrix} \mathrm{Id} \\ a-a(\delta,1) \end{pmatrix} \cdot \begin{pmatrix} \varphi \\ a(\delta,1) \end{pmatrix}$  and it has length  $<\delta$ .

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