

CHARACTERISTIC CLASSES FOR G -STRUCTURES

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ABSTRACT. Let $G \subset GL(V)$ be a linear Lie group with Lie algebra \mathfrak{g} and let $A(\mathfrak{g})^G$ be the subalgebra of G -invariant elements of the associative supercommutative algebra $A(\mathfrak{g}) = S(\mathfrak{g}^*) \otimes \Lambda(V^*)$. To any G -structure $\pi : P \rightarrow M$ with a connection ω we associate a homomorphism $\mu_\omega : A(\mathfrak{g})^G \rightarrow \Omega(M)$. The differential forms $\mu_\omega(f)$ for $f \in A(\mathfrak{g})^G$ which are associated to the G -structure π can be used to construct Lagrangians. If ω has no torsion the differential forms $\mu_\omega(f)$ are closed and define characteristic classes of a G -structure. The induced homomorphism $\mu'_\omega : A(\mathfrak{g})^G \rightarrow H^*(M)$ does not depend on the choice of the torsionfree connection ω and it is the natural generalization of the Chern Weil homomorphism.

1. G -STRUCTURES

1.1. G -structures. By a G -structure on a smooth finite dimensional manifold M we mean a principal fiber bundle $\pi : P \rightarrow M$ together with a representation $\rho : G \rightarrow GL(V)$ of the structure group in a real vector space V of dimension $\dim M$ and a 1-form σ (called the *soldering form*) on M with values in the associated bundle $P[V, \rho] = P \times_G V$ which is fiber wise an isomorphism and identifies $T_x M$ with $P[V]_x$ for each $x \in M$. Then σ corresponds uniquely to a G -equivariant 1-form $\theta \in \Omega_{\text{hor}}^1(P; V)^G$ which is strongly horizontal in the sense that its kernel is exactly the vertical bundle VP . The form θ is called the *displacement form* of the G -structure. A G -structure is called *1-integrable* if it admits torsionfree connections, see 1.4 below.

We fix this setting $((P, p, M, G), (V, \rho), \theta)$ from now on.

1.2. Invariant forms. We consider a multilinear form $f \in \bigotimes^k V^* = L^k(V)$ which is invariant in the sense that $f \circ (\bigotimes^k \rho(g)) = f$ for each $g \in G$. Let us denote by

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$L^k(V)^G$ the space of all these invariant forms. For each $f \in L^k(V)^G$ we have for any $X \in \mathfrak{g}$, the Lie algebra of G ,

$$\begin{aligned} 0 &= \frac{d}{dt} \Big|_0 f(\rho(\exp(tX))v_1, \dots, \rho(\exp(tX))v_k), \\ &= \sum_{i=1}^k f(v_1, \dots, \rho'(X)v_i, \dots, v_k), \end{aligned}$$

where $\rho' = T_e \rho : \mathfrak{g} \rightarrow \mathfrak{gl}(V)$ is the differential of the representation ρ .

1.3 Products of differential forms. For $\varphi \in \Omega^p(P; \mathfrak{g})$ and $\Psi \in \Omega^q(P; V)$ let us define the form $\rho'_\wedge(\varphi)\Psi \in \Omega^{p+q}(P; V)$ by

$$\begin{aligned} (\rho'_\wedge(\varphi)\Psi)(X_1, \dots, X_{p+q}) &= \\ &= \frac{1}{p!q!} \sum_{\sigma} \text{sign}(\sigma) \rho'(\varphi(X_{\sigma_1}, \dots, X_{\sigma_p})) \Psi(X_{\sigma_{p+1}}, \dots, X_{\sigma_{p+q}}). \end{aligned}$$

Then $\rho'_\wedge(\varphi) : \Omega^*(P; V) \rightarrow \Omega^{*+p}(P; V)$ is a graded $\Omega(P)$ -module homomorphism of degree p . Recall also that $\Omega(P; \mathfrak{g})$ is a graded Lie algebra with the bracket

$$\begin{aligned} [\varphi, \psi]_\wedge(X_1, \dots, X_{p+q}) &= \\ &= \frac{1}{p!q!} \sum_{\sigma} \text{sign} \sigma [\varphi(X_{\sigma_1}, \dots, X_{\sigma_p}), \psi(X_{\sigma_{p+1}}, \dots, X_{\sigma_{p+q}})]_{\mathfrak{g}}. \end{aligned}$$

One may easily check that for the graded commutator in $\text{End}(\Omega(P; V))$ we have

$$\rho'_\wedge([\varphi, \psi]_\wedge) = [\rho'_\wedge(\varphi), \rho'_\wedge(\psi)] = \rho'_\wedge(\varphi) \circ \rho'_\wedge(\psi) - (-1)^{pq} \rho'_\wedge(\psi) \circ \rho'_\wedge(\varphi)$$

so that $\rho'_\wedge : \Omega^*(P; \mathfrak{g}) \rightarrow \text{End}^*(\Omega(P; V))$ is a homomorphism of graded Lie algebras.

Let $\otimes V$ be the tensoralgebra generated by V . For $\Phi, \Psi \in \Omega(P; \otimes V)$ we will use the associative bigraded product

$$\begin{aligned} (\Phi \otimes_\wedge \Psi)(X_1, \dots, X_{p+q}) &= \\ &= \frac{1}{p!q!} \sum_{\sigma} \text{sign}(\sigma) \Phi(X_{\sigma_1}, \dots, X_{\sigma_p}) \otimes \Psi(X_{\sigma_{p+1}}, \dots, X_{\sigma_{p+q}}) \end{aligned}$$

1.4. The covariant exterior derivative. Let $\omega \in \Omega^1(P; \mathfrak{g})^G$ be a principal connection on the principal bundle (P, p, M, G) . Let $\chi : TP \rightarrow HP$ denote the corresponding projection onto the horizontal bundle $HP := \ker \omega$. The covariant exterior derivative $d_\omega : \Omega^k(P; V) \rightarrow \Omega_{\text{hor}}^{k+1}(P; V)$ is then given as usual by $d_\omega \Psi = \chi^* d\Psi = (d\Psi) \circ \Lambda^{k+1}(\chi)$.

Lemma. For $\Psi \in \Omega_{\text{hor}}(P; V)^G$ the covariant exterior derivative is given by $d_\omega \Psi = d\Psi + \rho'_\wedge(\omega)\Psi$.

Proof. If we insert one vertical vector field, say the fundamental vector field ζ_X for $X \in \mathfrak{g}$, into $d_\omega \Psi$, we get 0 by definition. For the right hand side we use $i_{\zeta_X} \Psi = 0$ and

$\mathcal{L}_{\zeta_X} \Psi = \frac{\partial}{\partial t} \Big|_0 (\text{Fl}_t^{\zeta_X})^* \Psi = \frac{\partial}{\partial t} \Big|_0 \Psi \circ \Lambda^p(r^{\exp tX}) = \frac{\partial}{\partial t} \Big|_0 \rho(\exp(-tX)) \Psi = -\rho'(X) \Psi$
to get

$$\begin{aligned} i_{\zeta_X} (d\Psi + \rho'_\wedge(\omega) \Psi) &= i_{\zeta_X} d\Psi + di_{\zeta_X} \Psi + \rho'_\wedge(i_{\zeta_X} \omega) \Psi - \rho'_\wedge(\omega) i_{\zeta_X} \Psi \\ &= \mathcal{L}_{\zeta_X} \Psi + \rho'_\wedge(X) \Psi = 0. \end{aligned}$$

Let now all vector fields ξ_i be horizontal, then we get

$$\begin{aligned} (d_\omega \Psi)(\xi_0, \dots, \xi_k) &= (\chi^* d\Psi)(\xi_0, \dots, \xi_k) = d\Psi(\xi_0, \dots, \xi_k), \\ (d\Psi + \rho'_\wedge(\omega) \Psi)(\xi_0, \dots, \xi_k) &= d\Psi(\xi_0, \dots, \xi_k). \quad \square \end{aligned}$$

1.5. Definition. If $\theta \in \Omega_{\text{hor}}^1(P; V)^G$ is the displacement form of a G -structure then the *torsion* of the connection ω with respect to the G -structure is $\tau := d_\omega \theta = d\theta + \rho'_\wedge(\omega)\theta$.

Recall that a G -structure is called 1-integrable if it admits a connection without torsion. This notion has also been investigated in [Kolář, Vadovičová] where it was called prolongable.

1.6. Chern-Weil forms. For differential forms $\psi_i \in \Omega^{p_i}(P; V)$ and $f \in L^k(V) = (\otimes^k V)^*$ we can construct the following differential forms:

$$\begin{aligned} \psi_1 \otimes_\wedge \dots \otimes_\wedge \psi_k &\in \Omega^{p_1 + \dots + p_k}(P; V \otimes \dots \otimes V), \\ f^{\psi_1, \dots, \psi_k} &:= f \circ (\psi_1 \otimes_\wedge \dots \otimes_\wedge \psi_k) \in \Omega^{p_1 + \dots + p_k}(P). \end{aligned}$$

The exterior derivative of the latter one is clearly given by

$$\begin{aligned} d(f \circ (\psi_1 \otimes_\wedge \dots \otimes_\wedge \psi_k)) &= f \circ d(\psi_1 \otimes_\wedge \dots \otimes_\wedge \psi_k) = \\ &= f \circ \left(\sum_{i=1}^k (-1)^{p_1 + \dots + p_{i-1}} \psi_1 \otimes_\wedge \dots \otimes_\wedge d\psi_i \otimes_\wedge \dots \otimes_\wedge \psi_k \right). \end{aligned}$$

We also set $f^\psi := f^{\psi, \dots, \psi} = \text{alt } f(\psi, \dots, \psi)$ for $\psi \in \Omega^p(P; V)$. Note that the form $f^{\psi_1, \dots, \psi_k}$ is G -invariant and horizontal if all $\psi_i \in \Omega_{\text{hor}}^{p_i}(P; V)^G$ and $f \in L^k(V)^G$. It is then the pullback of a form on M .

1.7. Lemma. Let $0 \neq \psi \in \Omega^p(P; V)$ and $f \in L^k(V)$. Then we have:

$$f^\psi \neq 0 \iff \begin{cases} \text{alt } f \neq 0, & \text{if } p \text{ is odd,} \\ \text{sym } f \neq 0, & \text{if } p \text{ is even,} \end{cases}$$

where *alt* and *sym* are the natural projections onto $\Lambda(V^*)$ and $S(V^*)$, respectively. \square

1.8. Lemma. If $f \in L^k(V)^G$ is invariant then we have

$$f \circ \left(\sum_{i=1}^k (-1)^{p_1 + \dots + p_{i-1}} \psi_1 \otimes_\wedge \dots \otimes_\wedge \rho'_\wedge(\omega) \psi_i \otimes_\wedge \dots \otimes_\wedge \psi_k \right) = 0.$$

Proof. This follows from the infinitesimal condition of invariance for f given in 1.2 by applying the alternator. \square

2. OBSTRUCTIONS TO 1-INTEGRABILITY OF G -STRUCTURES

2.1. Proposition. *Let $\pi : P \rightarrow M$ be a G -structure and let $f \in L^k(V)^G$ be an invariant tensor. For arbitrary G -equivariant horizontal V -valued forms $\psi_i \in \Omega_{\text{hor}}^{p_i}(P; V)^G$ we consider the $(p_1 + \dots + p_k)$ -form $f^{\psi_1, \dots, \psi_k}$ on M as above. If there is a connection ω for the G -structure π such that $d_\omega \psi_i = 0$ for all i , then the form $f^{\psi_1, \dots, \psi_k}$ is closed.*

Proof. We use $d_\omega \psi_i = d\psi_i + \rho'_\wedge(\omega)\psi_i$ from lemma 1.4, and lemma 1.8, to obtain

$$df^{\psi_1, \dots, \psi_k} = f \circ \left(\sum_{i=1}^k (-1)^{p_1 + \dots + p_{i-1}} \psi_1 \otimes_\wedge \dots \otimes_\wedge d_\omega \psi_i \otimes_\wedge \dots \otimes_\wedge \psi_k \right) = 0. \quad \square$$

2.2. Corollary. *1. For a G -structure $\pi : P \rightarrow M$ with displacement form θ we have a natural homomorphism of associative algebras*

$$\begin{aligned} \nu : \Lambda(V^*)^G &\rightarrow \Omega(M), \\ f &\mapsto f^\theta = f(\theta, \dots, \theta). \end{aligned}$$

2. If the G -structure is 1-integrable then the image of ν consists of closed forms and we get an induced homomorphism

$$\nu^* : \Lambda(V^*)^G \rightarrow H^*(M).$$

If M and G are compact then ν^ is injective.*

Proof. If the G -structure is 1-integrable then there is a connection ω with vanishing torsion $\tau = d_\omega \theta = 0$. Then the result follows from proposition 2.1.

If G is compact, any torsionfree connection ω for $\pi : P \rightarrow M$ is the Levi-Civita connection for some Riemannian metric. Any form f^θ , which is parallel with respect to ω , is harmonic and can thus not be exact for compact M . So ν^* is injective. \square

Problem. Is the homomorphism ν^* injective for compact M but noncompact G ?

2.3. Remark. Given a principal connection ω on P there is the induced covariant exterior derivative $\nabla : \Omega^p(M; P[V]) \rightarrow \Omega^{p+1}(M; P[V])$ on the associated vector bundle $P[V]$. The soldering form (see 1.1) $\sigma : TM \rightarrow P[V]$ is an isomorphism of vector bundles and we may consider the pull back covariant derivative $\sigma^* \nabla$ on TM . Next we consider the ‘combined’ covariant derivative $D^{\sigma^* \nabla, \nabla}$ on the vector bundle $L(TM, P[V])$ given by $D_X^{\sigma^* \nabla, \nabla} A = \nabla_X \circ A - A \circ (\sigma^* \nabla)_X$. Obviously we have $D^{\sigma^* \nabla, \nabla} \sigma = 0$. Consequently for any $f \in L^k(V)^G$ we have that $f^\theta \in \Omega^k(M)$ is parallel for the connection induced on $\Lambda^k T^* M$ from $\sigma^* \nabla$ on TM .

3. THE GENERALIZED CHERN-WEIL HOMOMORPHISM FOR G -STRUCTURES

3.1. The Chern-Weil homomorphism. Let ω be a connection for a G -structure $\pi : P \rightarrow M$ with curvature form $\Omega \in \Omega_{\text{hor}}^2(P, \mathfrak{g})$. Then the Bianchi identity $d_\omega \Omega = 0$ holds. If we apply proposition 2.1 to $\psi_i = \Omega$ we obtain a homomorphism

$$\gamma : S(\mathfrak{g}^*)^G \rightarrow \Omega(M),$$

given by $\gamma(f) = f^\Omega$. Since the image of γ consists of closed forms we have an induced homomorphism

$$\gamma' : S(\mathfrak{g}^*)^G \rightarrow H^*(M).$$

This is the well known Chern-Weil homomorphism.

3.2. The algebra $A(\mathfrak{g}, V)$. In order to generalize the Chern Weil homomorphism we associate to a Lie algebra \mathfrak{g} and a vector space V the associative graded commutative algebra

$$A(\mathfrak{g}, V) := S(\mathfrak{g}^*) \otimes \Lambda(V^*),$$

where the generators of the symmetric algebra $S(\mathfrak{g}^*)$ have degree 2. We may also consider $A(\mathfrak{g}, V)$ as a graded Lie algebra with the bracket

$$[a \otimes \varphi, b \otimes \psi] := \{a, b\} \otimes \varphi \wedge \psi, \quad a, b \in S(\mathfrak{g}^*), \varphi, \psi \in \Lambda(V^*),$$

where $\{a, b\}$ is the usual Poisson-Lie bracket in $S(\mathfrak{g}^*)$.

Let now \mathfrak{g} be the Lie algebra of the Lie group G and let $\rho : G \rightarrow GL(V)$ be a representation. Then G acts naturally on $A(\mathfrak{g}, V)$, and we denote $A(\mathfrak{g}, V)^G$ the subalgebra of G -invariant elements in $A(\mathfrak{g}, V)$.

3.3. Remark. The associative algebra $A(\mathfrak{g}, V)^G$ contains the subalgebra $S(\mathfrak{g}^*)^G \otimes \Lambda(V^*)^G$, in general as a proper subalgebra. Actually, let $G \subset GL(V)$ be the isotropy group of an irreducible Riemannian symmetric space M . Then the curvature tensor of M defines an element of $(\mathfrak{g}^* \otimes \Lambda^2 V^*)^G \subset A(\mathfrak{g}, V)^G$ that does not belong to $(\mathfrak{g}^*)^G \otimes (\Lambda^2 V^*)^G = 0$

3.4. The generalized Chern-Weil homomorphism. Now we are in a position to combine the constructions 2.2 and 3.1.

Theorem. *Let $\pi : P \rightarrow M$ be a G -structure on M with displacement form θ . Any connection ω in π defines a homomorphism of associative algebras*

$$\begin{aligned} \mu : A(\mathfrak{g}, V)^G &\rightarrow \Omega(M) \\ (S^p(\mathfrak{g}^*) \otimes \Lambda^q V^*)^G \ni f &\mapsto f^{\Omega, \theta} = f(\underbrace{\Omega, \dots, \Omega}_p, \underbrace{\theta, \dots, \theta}_q) \end{aligned}$$

If the connection ω has no torsion then the image of μ consists of closed forms and μ induces a homomorphism

$$\mu' : A(\mathfrak{g}, V)^G \rightarrow H^*(M),$$

which is independent of the choice of the torsionfree connection.

In other words, any G -invariant tensor $f \in S^p(\mathfrak{g}^*) \otimes \Lambda^q(V^*)$ defines a cohomology class $[f^{\Omega, \theta}] \in H^{2p+q}(M)$ which is an invariant of the 1-integrable G -structure. We call it a *characteristic class* of the 1-integrable G -structure π .

Proof. It just remains to show that the cohomology class $[f^{\Omega, \theta}]$ does not depend on the choice of the torsionfree connection for the G -structure $\pi : P \rightarrow M$.

So let ω_0, ω_1 be two torsionfree connections for the G -structure, let $\varphi = \omega_1 - \omega_0$, and denote by $\Omega_t = d_{\omega_t} \Omega_t$ the curvature form of the torsionfree connection $\omega_t = \omega_0 + t\varphi = (1-t)\omega_0 + t\omega_1$. We claim that for $f \in (S^p(\mathfrak{g}^*) \otimes \Lambda^q(V^*))^G$ we have

$$(1) \quad \begin{aligned} f^{\Omega_1, \theta} - f^{\Omega_0, \theta} &= d(Tf), \quad \text{where} \\ Tf &= p \int_0^1 f(\varphi, \Omega_t, \dots, \Omega_t, \theta, \dots, \theta) dt \end{aligned}$$

is the transgression form of f on P . The assertion is immediate from (1). To prove it we compute $\partial_t f^{\Omega_t, \theta}$ using the identities $\partial_t \Omega_t = d_{\omega_t} \varphi$ (see [Kobayashi, Nomizu II, p. 296]), $d_{\omega_t} \Omega_t = 0$, and $d_{\omega_t} \theta = 0$.

$$\begin{aligned} \partial_t f^{\Omega_t, \theta} &= p f(\partial_t \Omega_t, \Omega_t, \dots, \Omega_t, \theta, \dots, \theta) \\ &= p f(d_{\omega_t} \varphi, \Omega_t, \dots, \Omega_t, \theta, \dots, \theta) \\ &= p d_{\omega_t} f(\varphi, \Omega_t, \dots, \Omega_t, \theta, \dots, \theta) \\ &= p d f(\varphi, \Omega_t, \dots, \Omega_t, \theta, \dots, \theta). \quad \square \end{aligned}$$

3.5. Remarks about secondary characteristic classes. If the characteristic forms $f^{\Omega_1, \theta}$ and $f^{\Omega_0, \theta}$ associated with two torsionfree connections ω_1 and ω_0 vanish we obtain a secondary characteristic class $[Tf]$. It is a natural generalization of the classical Chern-Simons characteristic class, see [Chern, Simons], [Kobayashi, Ochiai].

Problem: study conditions when the secondary characteristic class $[Tf]$ does not depend on the choice of the torsionfree connections ω_1 and ω_0 .

3.6. Examples of characteristic classes. Assume that a linear group $G \subset GL(V)$ preserves some pseudo Euclidean metric in $V = \mathbb{R}^n$. Then we may identify the Lie algebra $\mathfrak{g} = \text{Lie}(G)$ with a subspace $\mathfrak{g} \subset \Lambda^2 V$. Suppose that there exists a G -invariant supplement \mathfrak{d} to \mathfrak{g} in $\Lambda^2 V$. Then the G -equivariant projection $\Lambda^2 V \rightarrow \mathfrak{g}$ along \mathfrak{d} determines a G -invariant element $q \in \mathfrak{g} \otimes \Lambda^2 V^* \cong \mathfrak{g}^* \otimes \Lambda^2 V^*$. The element q defines a 4-form $q^{\Omega, \theta}$ on the base of any G -structure $\pi : P \rightarrow M$ with a connection ω and curvature Ω . It may be written as

$$q^{\Omega, \theta} = q(\Omega, \theta, \theta) = q_{bcd}^a R_{aef}^b \theta^c \wedge \theta^d \wedge \theta^e \wedge \theta^f,$$

where (q_{bcd}^a) is the coordinate expression of q in the standard basis (e_a) of $V = \mathbb{R}^n$, $\theta = e_a \otimes \theta^a$, and $\Omega = R_{bef}^a \theta^e \wedge \theta^f$.

If ω is torsionfree the 4-form $q^{\Omega, \theta}$ is closed and it defines a cohomology class $[q^{\Omega, \theta}] \in H^4(M)$ independently of the choice of ω .

3.7. Remarks about the classification of characteristic classes. The classification of characteristic classes for G -structures with a given Lie group G reduces to the construction of generators of the associative algebra $A(\mathfrak{g}, V)^G = (S(\mathfrak{g}^*) \otimes \Lambda(V^*))^G$. We may also use the bracket to multiply characteristic classes. It suffices to solve this problem for those Lie groups G which appear as holonomy groups of torsionfree connection. Only for such groups G there exist 1-integrable non-flat G -structures. Under the additional hypothesis of irreducibility, all such groups were classified by [Berger], up to some gaps which were filled by [Bryant] and [Aleksievsky, Graev].

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